

# Chemical Science

Accepted Manuscript

This article can be cited before page numbers have been issued, to do this please use: J. Zhou, L. Stojanovic, A. A. Berezin, T. Battisti, A. Gill, B. Kariuki, D. Bonifazi, R. Crespo-Otero, M. R. Wasielewski and Y. Wu, *Chem. Sci.*, 2020, DOI: 10.1039/D0SC04646A.



This is an Accepted Manuscript, which has been through the Royal Society of Chemistry peer review process and has been accepted for publication.

Accepted Manuscripts are published online shortly after acceptance, before technical editing, formatting and proof reading. Using this free service, authors can make their results available to the community, in citable form, before we publish the edited article. We will replace this Accepted Manuscript with the edited and formatted Advance Article as soon as it is available.

You can find more information about Accepted Manuscripts in the [Information for Authors](#).

Please note that technical editing may introduce minor changes to the text and/or graphics, which may alter content. The journal's standard [Terms & Conditions](#) and the [Ethical guidelines](#) still apply. In no event shall the Royal Society of Chemistry be held responsible for any errors or omissions in this Accepted Manuscript or any consequences arising from the use of any information it contains.

# Organic room-temperature phosphorescence from halogen-bonded organic frameworks: hidden electronic effects in rigidified chromophores†

Jiawang Zhou,<sup>a</sup> Ljiljana Stojanović,<sup>b</sup> Andrey A. Berezin,<sup>c</sup> Tommaso Battisti,<sup>c</sup> Abigail Gill,<sup>c</sup> Benson M. Kariuki,<sup>c</sup> Davide Bonifazi,<sup>c,d</sup> Rachel Crespo-Otero,<sup>\*b</sup> Michael R. Wasielewski<sup>\*a</sup> and Yi-Lin Wu<sup>\*c</sup>

<sup>a</sup>Department of Chemistry and Institute for Sustainability and Energy at Northwestern, Northwestern University, Evanston, Illinois 60208-3113, United States. E-mail: m-wasielewski@northwestern.edu

<sup>b</sup>School of Biological and Chemical Sciences, Queen Mary University of London, London E1 4NS, United Kingdom. E-mail: r.crespo-otero@qmul.ac.uk

<sup>c</sup>School of Chemistry, Cardiff University, Cardiff CF10 3AT, United Kingdom. E-mail: WuYL@cardiff.ac.uk

<sup>d</sup>Institute of Organic Chemistry, Faculty of Chemistry, University of Vienna, Währinger Str. 38, 1090, Vienna, Austria.

† Electronic supplementary information (ESI) available: Experimental procedures, additional transient absorption data, synthetic and computational details, and X-ray crystallographic data. CCDC 1949875, 1949880 and 1949883. For ESI and crystallographic data in CIF or other electronic format see DOI: 10.1039/xxxxxxxxxx

**Abstract:** Development of purely organic materials displaying room-temperature phosphorescence (RTP) will expand the toolbox of inorganic phosphors for imaging, sensing or display applications. While molecular solids were found to suppress non-radiative energy dissipation and make the RTP process kinetically favourable, such an effect should be enhanced by the presence of multivalent directional non-covalent interactions. Here we report phosphorescence of a series of fast triplet-forming tetraethyl naphthalene-1,4,5,8-tetracarboxylates. Various numbers of bromo substituents were introduced to modulate intermolecular halogen-bonding interactions. Bright RTP with quantum yields up to 20% was observed when the molecule is surrounded by Br...O halogen-bonded network. Spectroscopic and computational analyses revealed that judicious heavy-atom positioning suppresses non-radiative relaxation and enhances intersystem crossing at the same time. The latter effect was found to be facilitated by the orbital angular momentum change, in addition to the conventional



*heavy-atom effect. Our results suggest the potential of multivalent non-covalent interactions for excited-state conformation and electronic control.*

View Article Online

DOI: 10.1039/D0SC04646A

## Introduction

Room temperature phosphorescence (RTP) has received increasing interest due to the potential it presents for photonic devices, bio-imaging, anti-counterfeiting, and night-vision applications.<sup>1-3</sup> Until recent years, the main sources of RTP luminophores have been inorganic or organometallic complexes, due to the presence of metal atoms being able to promote singlet-to-triplet intersystem crossing (ISC) in the excited states. However, heavy metal complexes or inorganic materials can often be toxic and expensive; through the study of purely organic phosphors, the applications of phosphorescence materials can expand by becoming more biocompatible, cheaper to acquire, and environmentally safer.<sup>4,5</sup> While there are many benefits of organic phosphors compared to those containing heavy metals, achieving RTP from purely organic molecules has proven a challenge on account of slow ISC rates and competitive non-radiative processes, in particular.

In recent decades, organic phosphorescence has become a more widely explored topic due to the discovery of long-lasting RTP by utilising crystallisation,<sup>6-8</sup> aggregation,<sup>9,10</sup> halogen bonding,<sup>11-14</sup> heavy atoms,<sup>15</sup> and carbonyl substituents<sup>16-18</sup> to circumvent the aforementioned issues.<sup>19-28</sup> Although spin-orbit coupling (SOC) in organic molecules is usually small (on the order of  $1\text{ cm}^{-1}$ , cf.  $10^2\text{--}10^3\text{ cm}^{-1}$  for organometallic complexes), the introduction of a carbonyl functionality to aromatic rings often opens up a  $^1(n-\pi^*) \rightarrow ^3(\pi-\pi^*)$  (or  $^1(\pi-\pi^*) \rightarrow ^3(n-\pi^*)$ ) channel with  $\text{SOC} \sim 100\text{ cm}^{-1}$ .<sup>29-33</sup> Such a small increase is sufficient to allow efficient ISC and populate the triplet of, for instance, benzophenone or benzaldehyde with a near-unitary quantum efficiency.<sup>34,35</sup> The structure of the as-generated triplet states can be rigidified in the solid state with the aid of non-covalent interactions (e.g. hydrogen and halogen bonds)<sup>11-13,20</sup> to suppress non-radiative vibrational relaxation, resulting in nearly quantitative RTP quantum yields in the solid state.<sup>19,36</sup>

Combining these design principles, the Kim group reported seminal work on efficient RTP luminophores based on 2,5-bis(hexyloxy)-4-bromobenzaldehyde.<sup>37</sup> The linear  $\text{C}=\text{O}\cdots\text{Br}$  halogen-bonding interactions<sup>38,39</sup> present in the solid state were suggested to be the major reason to avoid energy dissipation through vibrational motions. The proximity of a fourth-row Br element to the  $\text{C}=\text{O}$  group, where the non-bonding electrons originate in the  $n-\pi^*$  transition, is believed to enhance SOC as well.<sup>40,41</sup> In fact, in a later study by Kim and Dunietz, it was



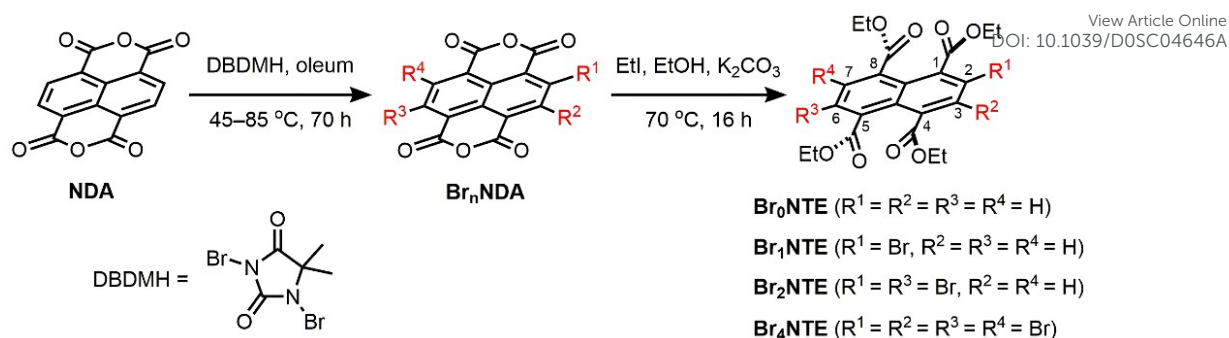
found that moving the Br substituent from the *para* to the *ortho* position, closer to the triplet-producing carbonyl functionality, in benzaldehyde increases SOC on the single-molecule level, which enhances both the rates of ISC  $k_{ISC}$  and phosphorescence  $k_{Phos}$  significantly by 5–15 fold.<sup>40</sup>

Inspired by these findings as well as other successful demonstration of halogen-bond-induced phosphorescence in the solid state, we exploited the naphthalene scaffold, a prototypical building block in organic optoelectronics, to study the effect of the halogen substitution and the role of halogen bonding in mediating triplet formation. Compared to the previously studied bromobenzaldehydes, this system is expected to have less carbonyl-originated  $n-\pi^*$  character in the low singlet excited states to drive ISC, thus offering a platform to highlight the halogen effects. Well-developed synthetic methodologies<sup>42–45</sup> were used to introduce multiple halogen-bond donors (e.g. Br) and acceptors (e.g. O) in naphthalene to permit multiple non-covalent interactions to occur synergistically, enabling phosphorescence from halogen-bonded frameworks.<sup>46</sup>

## Results and discussion

Naphthalene derivatives with halogen-bond accepting carbonyl functionalities and a various number of halogen-bond donating Br atoms can be prepared readily from 1,4,5,8-naphthalenetetracarboxylic dianhydride (NDA) (Scheme 1).<sup>47–49</sup> Bromination of NDA with 1,3-dibromo-5,5-dimethylhydantoin is slow and can produce a mixture of NDA with various numbers of Br substituents.<sup>47</sup> However, the application of excess reagents at elevated temperatures for a prolonged reaction time gives tetrabrominated NDA (**Br<sub>4</sub>NDA**) as the sole product. Esterification of **Br<sub>n</sub>NDA** with ethyl iodide in alkaline ethanol gave a mixture of naphthalene tetracarboxylic ethyl esters, **Br<sub>n</sub>NTE** ( $n = 0, 1, 2, 4$ ;  $n = 3$  can be isolated but it is not discussed here for simplicity).<sup>50, 51</sup> The individual compounds were isolated by column chromatography on SiO<sub>2</sub> and their identity was confirmed by NMR, MS, and single-crystal X-ray crystallography.

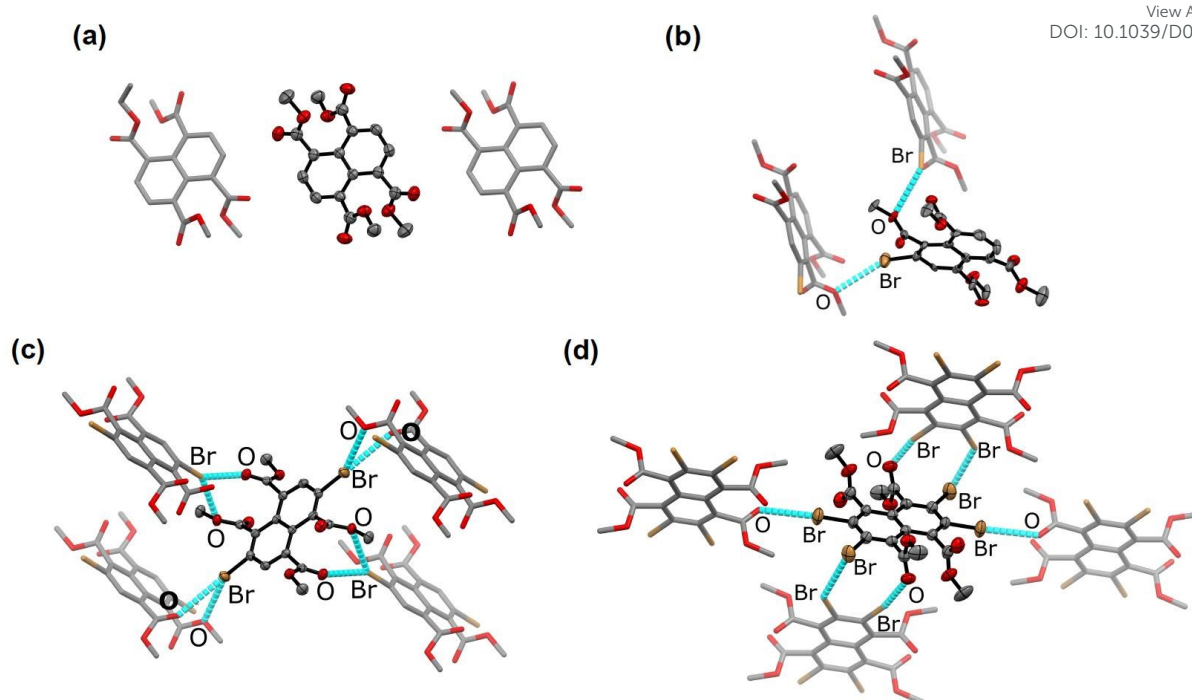




**Scheme 1** Synthesis of brominated naphthalene tetracarboxylic ethyl ester (**Br<sub>n</sub>NTE**).

In the solid state, an extended  $\text{Br}\cdots\text{O}$  network can be observed for **Br<sub>2</sub>NTE**<sup>50</sup> (Fig. 1). For each molecule, each pair of the *peri* ester groups interacts with a Br atom of a neighbouring molecule to establish bifurcated, slightly asymmetric halogen bonds with  $d(\text{C}-\text{Br}\cdots\text{O}=\text{C}) = 3.268(2)$  Å,  $d(\text{C}-\text{Br}\cdots\text{O}-\text{C}_2\text{H}_5) = 3.308(2)$  Å and both  $\theta(\text{C}-\text{Br}\cdots\text{O}) \sim 150^\circ$  (*i.e.*  $152.40(9)$  and  $149.21(8)$ ).<sup>38</sup> Reciprocally, each Br atom is interacting with two *peri* ester groups of a nearby molecule of **Br<sub>2</sub>NTE**. Being symmetrically substituted with four Br and four ester functionalities, **Br<sub>4</sub>NTE** is also embedded in a framework of halogen bonds in the solid state. However, likely due to the steric requirement of large Br atoms, the same arrangement in **Br<sub>2</sub>NTE** was not observed. Instead, only two out of the four esters on the 1 and 5 positions form linear  $\text{C}-\text{Br}\cdots\text{O}=\text{C}$  short contacts with the Br atoms on the 3 and 7 positions of the neighbouring molecules (values taken from two crystallographically independent molecules:  $d(\text{Br}\cdots\text{O}) = 3.074(3)$  and  $3.286(3)$  Å,  $\theta(\text{C}-\text{Br}\cdots\text{O}) = 165.5(1)$  and  $168.5(1)^\circ$ ). The remaining two Br atoms on the 2 and 6 positions are engaged in orthogonal (“Type II”)<sup>52, 53</sup>  $\text{C}-\text{Br}\cdots\text{Br}-\text{C}$  interactions ( $d(\text{Br}\cdots\text{Br}) = 3.692(1)$  Å,  $\theta(\text{C}-\text{Br}\cdots\text{Br}) = 86.4(1)^\circ$ ).



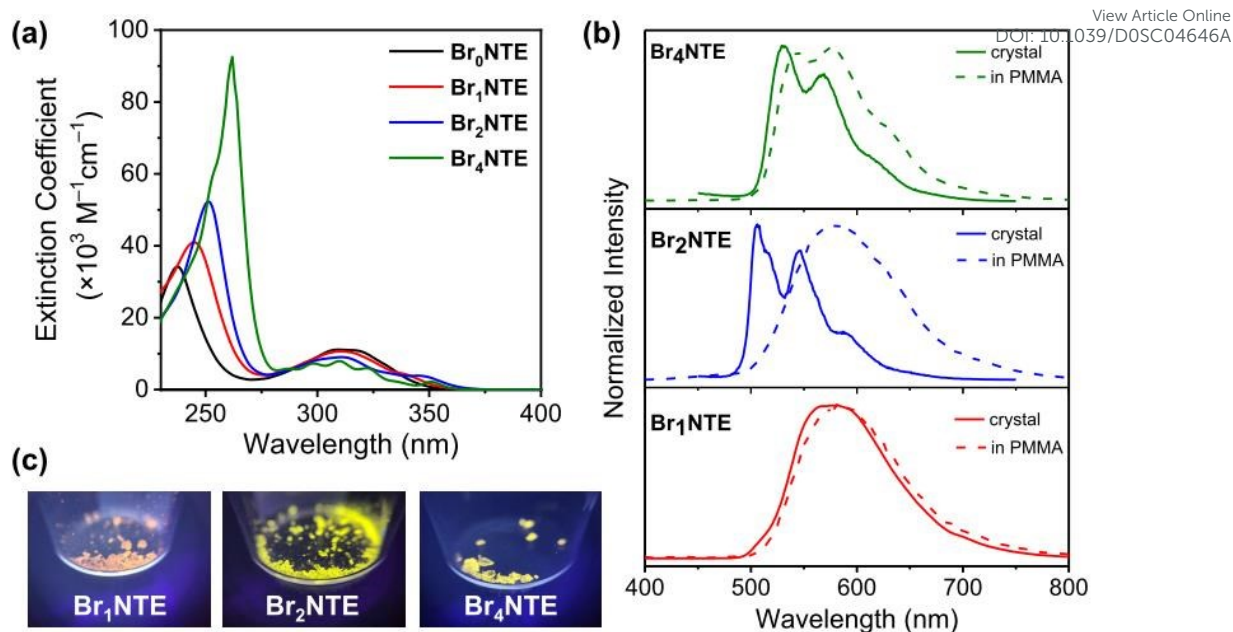


**Fig. 1** Single crystal X-ray molecular structure of (a) **Br<sub>0</sub>NTE** (space group  $P2_1/n$ ), (b) **Br<sub>1</sub>NTE** ( $P2_1$ ), (c) **Br<sub>2</sub>NTE** ( $P2_1/n$ ), and (d) **Br<sub>4</sub>NTE** ( $P-1$ ) and the close neighbours in crystals. Crystals were obtained by diffusing MeOH vapour into the  $\text{CHCl}_3$  solutions of **Br<sub>n</sub>NTE**. Colour code: C = grey, O = red, Br = brown. For **Br<sub>1</sub>NTE**, only the major component of the disorder is shown and discussed. The terminal carbon of the ethyl group and all hydrogen atoms are omitted for clarity. Thermal ellipsoids of the central molecules are shown at the 50% probability level, whereas the neighbouring molecules shown in stick representation. Non-covalent  $\text{Br}\cdots\text{O}$  and  $\text{Br}\cdots\text{Br}$  interactions are highlighted with cyan dashed lines.

With only one Br atom per molecule, the ester groups in **Br<sub>1</sub>NTE** do not engage in extended halogen-bonded networks. In fact, the shortest  $d(\text{C}-\text{Br}\cdots\text{O}-\text{C}_2\text{H}_5)$  distance is measured to be  $3.32(2)$  Å ( $\theta(\text{C}-\text{Br}\cdots\text{O}) = 151.7(9)^\circ$ ), barely shorter than the van der Waals contact distance. At last, no  $\pi$ -stack or short contact between C–H and naphthalene was found in the crystals of **Br<sub>0</sub>NTE**.







**Fig. 2** (a) UV-Vis absorption spectra of  $\text{Br}_n\text{NTE}$  at  $(5-8) \times 10^{-5} \text{ M}$  in  $\text{CH}_2\text{Cl}_2$ . (b) Normalised phosphorescence spectra of  $\text{Br}_n\text{NTE}$  in the crystalline solid state (solid line) or in PMMA (dashed line). Samples were excited at 300–320 nm. (c) Photographs of solid emission under UV irradiation (365 nm).

While naphthalene and its 2-brominated derivatives display electronic absorption  $< 300 \text{ nm}$ , ester substitution induces a bathochromic shift of the naphthalene-centred transitions by *ca.* 50 nm, extending the absorption bands to 350 nm with maxima at  $\sim 300 \text{ nm}$  (Fig. 2).<sup>54-56</sup> Compared to pristine naphthalene, which has an appreciable fluorescence quantum yield of 23% (40% triplet formation yield),<sup>55</sup> no emission was detected from all  $\text{Br}_n\text{NTE}$  in deaerated  $\text{CH}_2\text{Cl}_2$  up to 0.02 M (near saturation) excited at 330 nm.

Despite the non-radiative energy dissipation in solution, crystalline solids of the brominated molecules display visible phosphorescence in the 500–700 nm region with millisecond lifetimes, whereas non-brominated  $\text{Br}_0\text{NTE}$  remains non-emissive (Fig. 2 and Table 1). Powdery crystalline solid samples of  $\text{Br}_n\text{NTE}$ , whose powder X-ray diffraction profiles match the pattern based on their single-crystal data, were used in all phosphorescence measurements. Phosphorescence of crystalline  $\text{Br}_2\text{NTE}$  and  $\text{Br}_4\text{NTE}$  feature clear vibrational progression with a quantum yield of  $\Phi_{\text{Phos}} = 19.6\%$  and  $9.3\%$ , respectively. Much weaker and structureless emission was observed for  $\text{Br}_1\text{NTE}$  ( $\Phi_{\text{Phos}} = 1.4\%$ ).

The varying luminescent behaviours suggest that the excited-state dynamics were modulated in a subtle way by Br-specific properties, which is however not directly related to the number of Br atoms in the molecule. It is conceivable that multi-point halogen bonding provides a geometric framework to strengthen the rigidity of  $\text{Br}_n\text{NTE}$  in the crystalline state.



This effect is especially substantial for **Br<sub>2</sub>NTE** where all the peripheral substituents engage in the directional Br $\cdots$ O interactions, providing the additional factor to the solid-state effect<sup>6, 8</sup> of RTP to impede competitive non-radiative relaxation through intramolecular motions. The highest phosphorescence quantum yield was thus observed for the crystalline sample of **Br<sub>2</sub>NTE**.

The weaker and non-structured phosphorescence observed for **Br<sub>1</sub>NTE** (and **Br<sub>0</sub>NTE**) seems to be originated from its looser solid-state packing. If we define the volumetric index  $V_i$  as the ratio between the Voronoi volume ( $V_{\text{Vor}}$ )<sup>57, 58</sup> and the van der Waals volume ( $V_{\text{vdW}}$ ) of a molecule in the crystal, smaller  $V_i = V_{\text{Vor}} / V_{\text{vdW}}$  would suggest denser packing.  $V_i$  of 1.27–1.30 were found for **Br<sub>2</sub>NTE** and **Br<sub>4</sub>NTE** embedded in halogen-bonded frameworks, but the values are significantly larger for **Br<sub>0</sub>NTE** and **Br<sub>1</sub>NTE** (1.36–1.38). The larger free space available to each molecule in the **Br<sub>0</sub>NTE** and **Br<sub>1</sub>NTE** crystals allows the excited molecules to decay radiatively and non-radiatively on various points of the triplet potential energy surface.

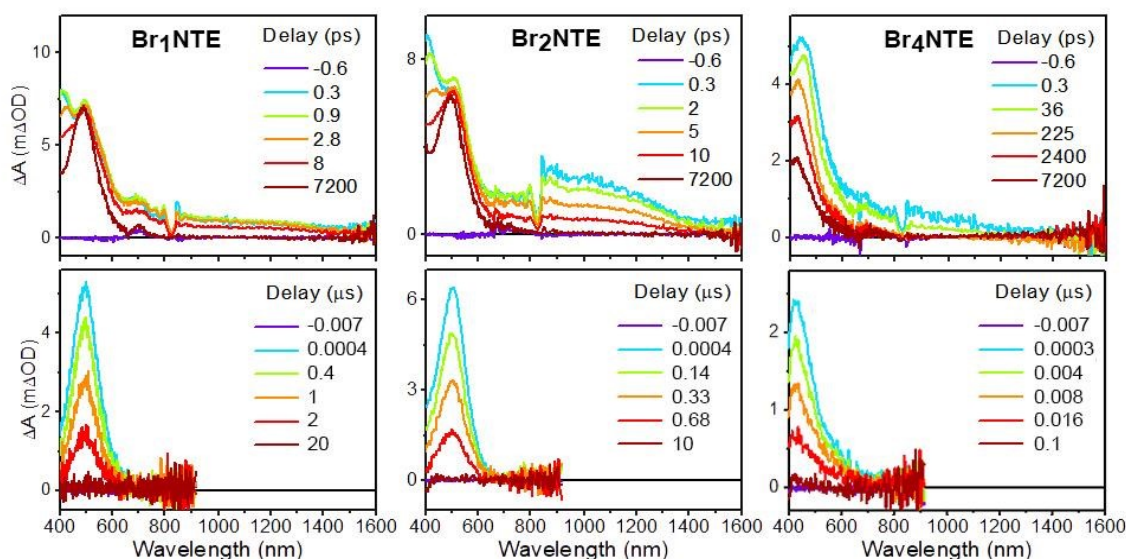
The significance of the inter-**Br<sub>n</sub>NTE** Br $\cdots$ O interactions is further supported by comparing the phosphorescence of crystalline **Br<sub>n</sub>NTE** with that of the dispersed molecules in poly(methyl methacrylate) (PMMA, Mw ~ 996 kDa; 2 wt% doping). The rigid polymer matrix is expected to constrain the molecular motion at room temperature but disrupt inter-**Br<sub>n</sub>NTE** Br $\cdots$ O halogen bonds. The phosphorescence spectra of **Br<sub>1</sub>NTE** remained identical in either environment (Fig. 2), indicating that the triplet decay in **Br<sub>1</sub>NTE** is largely intrinsic to the monomeric molecule. However, the vibrational progression of **Br<sub>4</sub>NTE**, a signature of chromophore rigidity, became less pronounced, and that of **Br<sub>2</sub>NTE** completely disappeared and the overall emission profile resembles very well to that of **Br<sub>1</sub>NTE**.

Additional support for the efficient population of the triplet excited state was provided by transient absorption measurements. Spectroscopically, all **Br<sub>n</sub>NTE** exhibit similar excited-state dynamics: following the initial formation of the singlet excited state, which displays excited-state absorption (ESA) peaking at *ca.* 490 nm and a broad feature in the near infrared region of 800–1000 nm (Fig. 3 and ESI† Section 5), a new excited-state species with ESA at *ca.* 480 nm appears with microsecond lifetimes. This long-lived species was assigned to the triplet of each chromophore based on the lifetime and spectral similarity to the triplet-triplet absorption of methyl 1-naphthalate<sup>59</sup> and 2-bromonaphthalene.<sup>60</sup> Therefore, the decay of the initial state can be ascribed to singlet-to-triplet ISC; time constants on the order of tens of picosecond were observed for this process (Table 1). Compared to the typical fluorescence lifetime (1 ns or longer) of naphthalene derivatives,<sup>54, 61</sup> the fast ISC process suggests a high





triplet forming efficiency. Such efficient ISC on the molecular level is likely due to the combined results of bromo (*cf.* > 90% triplet yield for 2-bromonaphthalene)<sup>60</sup> and carbonyl substitution.<sup>62</sup> Broadly speaking, the more bromo atoms in a molecule, the faster the  $S_1 \rightarrow T_n$  and  $T_1 \rightarrow S_0$  processes, in line with the stronger heavy-atom enhanced SOC.<sup>63</sup> Unexpectedly, however, the  $S_1 \rightarrow T_n$  ISC for **Br<sub>4</sub>NTE** is noticeably slower than its less brominated analogues.



**Fig. 3** Transient absorption spectra of **Br<sub>n</sub>NTE** ( $n = 1, 2$ , and  $4$ ) in deaerated  $\text{CH}_2\text{Cl}_2$  at various pump-probe delay indicated (excitation = 330 nm, see ESI† Section 5 for  $n = 0$ ).

**Table 1.** Photophysical properties of **Br<sub>n</sub>NTE**.

	in $\text{CH}_2\text{Cl}_2$		crystalline solids	
	$\tau(S_1 \rightarrow T_n)^a$	$\tau(T_1 \rightarrow S_0)^a$	$\tau_{\text{Phos}}^b$	$\Phi_{\text{Phos}}^{b,c}$
<b>Br<sub>0</sub>NTE</b>	111.7 ± 0.8 ps <sup>d</sup>	29.1 ± 0.1 μs	n.a.	n.a.
<b>Br<sub>1</sub>NTE</b>	9.1 ± 0.3 ps <sup>d</sup>	1.42 ± 0.01 μs	1.53 ± 0.02 ms	1.4%
<b>Br<sub>2</sub>NTE</b>	7.5 ± 0.3 ps	0.50 ± 0.01 μs	1.94 ± 0.01 ms	19.6%
<b>Br<sub>4</sub>NTE</b>	48.3 ± 0.9 ps <sup>d</sup>	0.0121 ± 0.0006 μs	1.11 ± 0.01 ms	9.3%

<sup>a</sup> From transient absorption measurements. <sup>b</sup> From (time-resolved) phosphorescence measurements. <sup>c</sup> The uncertainty is estimated to be 20% of the measured values. <sup>d</sup> Preceded by the relaxation of hot  $S_1$  in  $(0.9\text{--}1.2) \pm 0.3$  ps.

Since the rate of  $S_1 \rightarrow T_n$  ISC is largely determined by the energy gap between the singlet and triplet states ( $\Delta E_{\text{ST}}$ ) and the magnitude of spin-orbit coupling (SOC),<sup>64</sup> we evaluated the matrix elements of  $\langle S_1 | \hat{H}_{\text{SO}} | T_n \rangle$  using the two-layer ONIOM (QM:MM) scheme to simulate the photophysical processes in crystals. The molecular geometry was computed at the level of  $\omega\text{B97X-D/6-31G(d):OPLS-AA}$ , and  $\langle S_1 | \hat{H}_{\text{SO}} | T_n \rangle$  calculated at the TDA- $\omega\text{B97X-D/6-}$



311+G(d,p) level of theory based on the ONIOM geometries (Table 2 and ESI† Section 7). The (TD-)DFT calculations were performed using Gaussian 16,<sup>65</sup> which was then interfaced with PySOC<sup>30</sup> to evaluate the SOC matrix elements. The Tamm–Dancoff approximation (TDA) was exploited to minimise triplet instability.<sup>40, 66</sup> In all cases, the state energies are not significantly affected by aggregation; thus results from the calculations with one molecule in the QM region are discussed here.

The ISC process for **Br<sub>n</sub>NTE** likely takes place between S<sub>1</sub> and the high-lying triplet states. Considering  $\Delta E_{ST}$  alone (<0.5 eV), ISC to T<sub>2,3</sub> for **Br<sub>1</sub>NTE**, T<sub>2-4</sub> for **Br<sub>2</sub>NTE**, and T<sub>2</sub> for **Br<sub>4</sub>NTE** should dominate in the respective molecules, whereas the large energy gap  $\Delta E_{ST}$  > 1.5 eV prevents direct ISC into T<sub>1</sub> (see ESI† Section 7 for the relative energies). Compared to **Br<sub>0</sub>NTE**, incorporating fourth-row Br elements into the naphthalene scaffold increases SOC by 1–2 orders of magnitude. Despite the larger number of Br atoms in the structure, smaller SOC was found for **Br<sub>4</sub>NTE** than **Br<sub>2</sub>NTE**, in line with the slower triplet formation found experimentally for the former molecule. The  $\langle S_0 | \hat{H}_{SO} | T_1 \rangle$  calculated at the T<sub>1</sub> geometry, the key factor determining the rate of phosphorescence, was similarly found to be smaller for **Br<sub>4</sub>NTE** than **Br<sub>2</sub>NTE**.

**Table 2.** Spin-orbit coupling (in cm<sup>-1</sup>) calculated at the TDA- $\omega$ B97X-D/6-311+G(d,p) level of theory based on the ONIOM geometries.

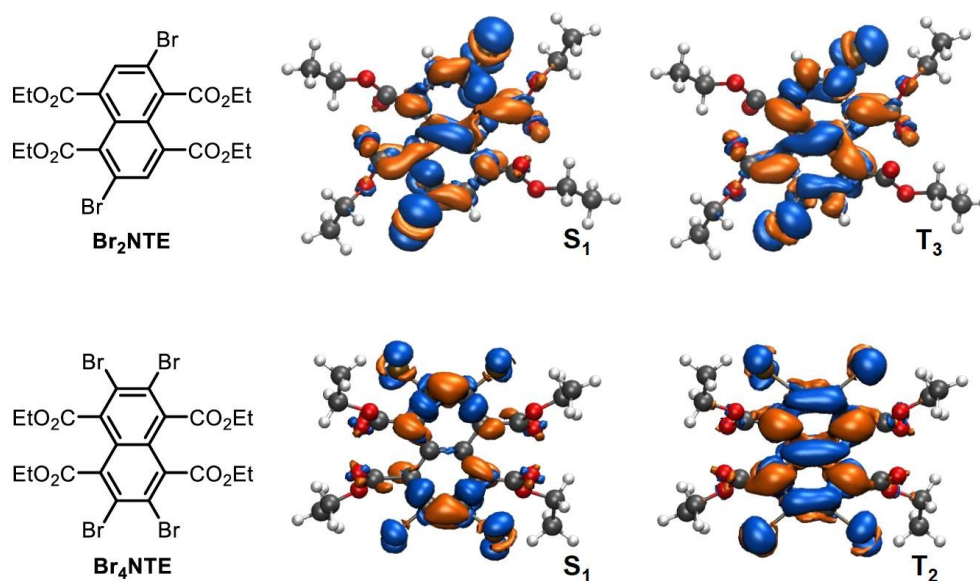
	at S <sub>1</sub> geometry <sup>a</sup>				at T <sub>1</sub> geometry
	$\langle S_1   \hat{H}_{SO}   T_1 \rangle$	$\langle S_1   \hat{H}_{SO}   T_2 \rangle$	$\langle S_1   \hat{H}_{SO}   T_3 \rangle$	$\langle S_1   \hat{H}_{SO}   T_4 \rangle$	$\langle S_0   \hat{H}_{SO}   T_1 \rangle$
<b>Br<sub>0</sub>NTE</b>	0.82	0.04	0.52	0.86	0.01
<b>Br<sub>1</sub>NTE</b>	8.22	<b>10.87</b>	<b>15.49</b>	9.02	3.22
<b>Br<sub>2</sub>NTE</b>	68.74	0.93	<b>166.79</b>	<b>106.08</b>	142.3
<b>Br<sub>4</sub>NTE</b>	20.30	<b>42.79</b>	3.55	30.93	0.38

<sup>a</sup> States relevant for the intersystem-crossing mechanism are highlighted in bold.

A close examination of the electron density of the key states provided hints to the origin of the unexpected drop in SOC for **Br<sub>4</sub>NTE**. Fig. 4 shows the electron density difference between the selected excited states and the ground state for **Br<sub>2</sub>NTE** (S<sub>1</sub> and T<sub>3</sub>) and for **Br<sub>4</sub>NTE** (S<sub>1</sub> and T<sub>2</sub>). These transitions displayed a significant naphthalene-centred  $\pi$ - $\pi^*$  character; the involvement of the Br atoms can be clearly seen and hence the higher SOC in brominated **Br<sub>n</sub>NTE**. Comparatively, the carbonyl n- $\pi^*$  contribution, the typical driver for the ISC process in aromatic ketones/aldehydes, appears to be much less substantial. In the case of **Br<sub>2</sub>NTE**, the Br-centred transition densities are roughly perpendicular to the naphthalene plane



in the  $S_1$  state but rotate distinctively in the  $T_3$  state, facilitating the orbital angular momentum change for ISC (similar rotation found in  $T_4$ ). In the case of **Br<sub>4</sub>NTE**, however, the Br-centred transition densities in  $S_1$  and  $T_2$  are both perpendicular to the naphthalene plane. The absence of the analogous rotated transition density for **Br<sub>4</sub>NTE** is understandable as unfavourable electron repulsion in the region between neighbouring Br atoms would be caused by such a change in density orientation.



**Fig. 4** Top-down view of electron density difference plots ( $0.001 \text{ e Bohr}^{-3}$  isovalue) between the selected excited states ( $S_1$  or  $T_{2/3}$ ) and the ground state for **Br<sub>2</sub>NTE** (top row) and **Br<sub>4</sub>NTE** (bottom row). The molecular orientation is sketched on the left; orange colour represents positive and blue negative values.

Taken together, the judicious heavy-atom positioning in **Br<sub>2</sub>NTE** results in the favourable structural and electronic contributions to its efficient RTP. The 2,6-dibromo substitution offers a lock-in mechanism through halogen bonding to inhibit non-radiative relaxation. Furthermore, high SOC and hence efficient ISC are made possible by adding the orbital angular momentum change to the heavy-atom effect in both the triplet-generation ( $S_1 \rightarrow T_n$ ) and phosphorescence ( $T_1 \rightarrow S_0$ ) processes.

## Conclusions

In summary, we have shown that simultaneously incorporating multiple heavy halogens and halogen-bond donor/acceptor pairs in aromatic molecules can enable bright phosphorescence from purely organic materials. The formation of halogen-bonded frameworks in the solid states rigidifies phosphorophores, favouring the radiative decay. However, our results indicate that a



fine balance has to be struck in terms of the number and positioning of halogens. Too many large halogen atoms in proximity may prohibit *structurally* the access of halogen-bond acceptors and *electronically* the contribution of the non-bonding electrons of halogens for enhancing SOC. The latter effect is especially important to consider in the case of carbonyl-bearing polycyclic aromatic hydrocarbons, such as rylenes and its derivatives in the present study where the  $S_1$  state is primarily  $\pi$ - $\pi^*$  in nature. It should be noted that the formation of halogen bonds cannot necessarily be correlated to the increase in ISC and phosphorescence rates; an excited-state analysis will be needed to elucidate the magnitude and origin of SOC when designing organic RTP materials.

## Conflicts of interest

There are no conflicts to declare.

## Acknowledgements

D.B. gratefully acknowledges the EU through the MSCA-ITN-ETN (project PHOTOTRAIN) and School of Chemistry at Cardiff University for generous financial support. R.C.-O. acknowledges funding from the EPSRC (EP/R029385/1) and Leverhulme Trust (RPG-2019-122). This research utilised Queen Mary's Apocrita HPC facility. This work was supported by the US Department of Energy, Office of Science, Office of Basic Energy Sciences under Award DE-SC0020168 (M.R.W.). The authors thank Professor Kenneth D. M. Harris and Dr Colan E. Hughes (Cardiff) for the help with powder X-ray diffraction.

## References

1. Z. W. Pan, Y. Y. Lu and F. Liu, *Nat. Mater.*, 2012, **11**, 58-63.
2. T. Maldiney, A. Bessiere, J. Seguin, E. Teston, S. K. Sharma, B. Viana, A. J. J. Bos, P. Dorenbos, M. Bessodes, D. Gourier, D. Scherman and C. Richard, *Nat. Mater.*, 2014, **13**, 418-426.
3. Y. Li, M. Gecevicius and J. R. Qiu, *Chem. Soc. Rev.*, 2016, **45**, 2090-2136.
4. S. M. A. Fateminia, Z. Mao, S. D. Xu, Z. Y. Yang, Z. G. Chi and B. Liu, *Angew. Chem. Int. Ed.*, 2017, **56**, 12160-12164.
5. G. Q. Zhang, G. M. Palmer, M. Dewhirst and C. L. Fraser, *Nat. Mater.*, 2009, **8**, 747-751.



6. Y. Gong, L. Zhao, Q. Peng, D. Fan, W. Z. Yuan, Y. Zhang and B. Z. Tang, *Chem. Sci.*, 2015, **6**, 4438-4444. View Article Online  
DOI: 10.1039/C5SC04646A
7. R. Yoshii, A. Hirose, K. Tanaka and Y. Chujo, *J. Am. Chem. Soc.*, 2014, **136**, 18131-18139.
8. W. Z. Yuan, X. Y. Shen, H. Zhao, J. W. Y. Lam, L. Tang, P. Lu, C. L. Wang, Y. Liu, Z. M. Wang, Q. Zheng, J. Z. Sun, Y. G. Ma and B. Z. Tang, *J. Phys. Chem. C*, 2010, **114**, 6090-6099.
9. Z. Zhao, X. Zheng, L. Du, Y. Xiong, et al., *Nat. Commun.*, 2019, **10**, 2952.
10. G. Bergamini, A. Fermi, C. Botta, U. Giovanella, S. Di Motta, F. Negri, R. Peresutti, M. Gingras and P. Ceroni, *J. Mater. Chem. C*, 2013, **1**, 2717-2724.
11. L. Xiao, Y. Wu, Z. Yu, Z. Xu, J. Li, Y. Liu, J. Yao and H. Fu, *Chem. Eur. J.*, 2018, **24**, 1801-1805.
12. S. K. Maity, S. Bera, A. Paikar, A. Pramanik and D. Halder, *Chem. Commun.*, 2013, **49**, 9051-9053.
13. S. Cai, H. Shi, D. Tian, H. Ma, Z. Cheng, Q. Wu, M. Gu, L. Huang, Z. An, Q. Peng and W. Huang, *Adv. Funct. Mater.*, 2018, **28**, 1705045.
14. Z. Yang, C. Xu, W. Li, Z. Mao, X. Ge, Q. Huang, H. Deng, J. Zhao, F. L. Gu, Y. Zhang and Z. Chi, *Angew. Chem. Int. Ed.*, 2020, DOI: 10.1002/anie.202007343.
15. A. Kremer, C. Aurisicchio, F. De Leo, B. Ventura, J. Wouters, N. Armaroli, A. Barbieri and D. Bonifazi, *Chem. Eur. J.*, 2015, **21**, 15377-15387.
16. J. Xu, A. Takai, Y. Kobayashi and M. Takeuchi, *Chem. Commun.*, 2013, **49**, 8447-8449.
17. D. Lee, O. Bolton, B. C. Kim, J. H. Youk, S. Takayama and J. Kim, *J. Am. Chem. Soc.*, 2013, **135**, 6325-6329.
18. M. Shimizu, A. Kimura and H. Sakaguchi, *Eur. J. Org. Chem.*, 2016, 467-473.
19. S. Mukherjee and P. Thilagar, *Chem. Commun.*, 2015, **51**, 10988-11003.
20. L. Xiao and H. Fu, *Chem. Eur. J.*, 2019, **25**, 714-723.
21. Z. He, W. Zhao, J. W. Y. Lam, Q. Peng, H. Ma, G. Liang, Z. Shuai and B. Z. Tang, *Nat. Commun.*, 2017, **8**, 416.
22. W. Zhao, Z. He, Jacky W. Y. Lam, Q. Peng, H. Ma, Z. Shuai, G. Bai, J. Hao and Ben Z. Tang, *Chem*, 2016, **1**, 592-602.
23. A. Barbieri, E. Bandini, F. Monti, V. K. Praveen and N. Armaroli, *Top. Curr. Chem.*, 2016, **374**, 47.
24. M. Hayduk, S. Riebe and J. Voskuhl, *Chem. Eur. J.*, 2018, **24**, 12221-12230.





25. M. S. Kwon, D. Lee, S. Seo, J. Jung and J. Kim, *Angew. Chem. Int. Ed.*, 2014, **53**, 11177-11181. View Article Online  
DOI: 10.1039/D0SC04646A
26. L. Bian, H. Shi, X. Wang, K. Ling, H. Ma, M. Li, Z. Cheng, C. Ma, S. Cai, Q. Wu, N. Gan, X. Xu, Z. An and W. Huang, *J. Am. Chem. Soc.*, 2018, **140**, 10734-10739.
27. M. Baroncini, G. Bergamini and P. Ceroni, *Chem. Commun.*, 2017, **53**, 2081-2093.
28. A. Forni, E. Lucenti, C. Botta and E. Cariati, *J. Mater. Chem. C*, 2018, **6**, 4603-4626.
29. R. Liu, X. Gao, M. Barbatti, J. Jiang and G. Z. Zhang, *J. Phys. Chem. Lett.*, 2019, **10**, 1388-1393.
30. X. Gao, S. M. Bai, D. Fazzi, T. Niehaus, M. Barbatti and W. Thiel, *J. Chem. Theory Comput.*, 2017, **13**, 515-524.
31. M. A. El-Sayed, *J. Chem. Phys.*, 1963, **38**, 2834-2838.
32. F. Dinkelbach, M. Kleinschmidt and C. M. Marian, *J. Chem. Theory Comput.*, 2017, **13**, 749-766.
33. D. Sasikumar, A. T. John, J. Sunny and M. Hariharan, *Chem. Soc. Rev.*, 2020, DOI: 10.1039/D0CS00484G.
34. R. H. Compton, K. T. V. Grattan and T. Morrow, *J. Photochem.*, 1980, **14**, 61-66.
35. T. Itoh, *Chem. Phys. Lett.*, 1988, **151**, 166-168.
36. S. Hirata, *Adv. Opt. Mater.*, 2017, **5**, 1700116.
37. O. Bolton, K. Lee, H. J. Kim, K. Y. Lin and J. Kim, *Nat. Chem.*, 2011, **3**, 205-210.
38. G. Cavallo, P. Metrangolo, R. Milani, T. Pilati, A. Priimagi, G. Resnati and G. Terraneo, *Chem. Rev.*, 2016, **116**, 2478-2601.
39. P. J. Costa, *Phys. Sci. Rev.*, 2017, **2**, 20170136.
40. S. Sarkar, H. P. Hendrickson, D. Lee, F. DeVine, J. Jung, E. Geva, J. Kim and B. D. Dunietz, *J. Phys. Chem. C*, 2017, **121**, 3771-3777.
41. S. J. Ang, T. S. Chwee and M. W. Wong, *J. Phys. Chem. C*, 2018, **122**, 12441-12447.
42. M. Al Kobaisi, S. V. Bhosale, K. Latham, A. M. Raynor and S. V. Bhosale, *Chem. Rev.*, 2016, **116**, 11685-11796.
43. O. Cakmak, *J. Chem. Res.*, 1999, 366-367.
44. A. M. Wagner, A. J. Hickman and M. S. Sanford, *J. Am. Chem. Soc.*, 2013, **135**, 15710-15713.
45. K. D. Collins, R. Honeker, S. Vasquez-Cespedes, D. T. D. Tang and F. Glorius, *Chem. Sci.*, 2015, **6**, 1816-1824.





46. N. Chongboriboon, K. Samakun, T. Inprasit, F. Kielar, W. Dungkaew, L. W. Y. Wong, H. H. Y. Sung, D. B. Ninkovic, S. D. Zaric and K. Chainok, *Crystengcomm*, 2020, **22**, 24-34. View Article Online  
DOI: 10.1039/D9CC04646A
47. A. A. Berezin, A. Sciutto, N. Demitri and D. Bonifazi, *Org. Lett.*, 2015, **17**, 1870-1873.
48. C. Roger and F. Wurthner, *J. Org. Chem.*, 2007, **72**, 8070-8075.
49. M. Sasikumar, Y. V. Suseela and T. Govindaraju, *Asian J. Org. Chem.*, 2013, **2**, 779-785.
50. R. S. K. Kishore, V. Ravikumar, G. Bernardinelli, N. Sakai and S. Matile, *J. Org. Chem.*, 2008, **73**, 738-740.
51. Y. Ma, X. Zhang, S. Stappert, Z. Yuan, C. Li and K. Müllen, *Chem. Commun.*, 2017, **53**, 5310-5313.
52. B. K. Saha, A. Saha and S. A. Rather, *Cryst. Growth Des.*, 2017, **17**, 2314-2318.
53. B. K. Saha, S. A. Rather and A. Saha, *Cryst. Growth Des.*, 2016, **16**, 3059-3062.
54. M. Kitamura and H. Baba, *Bull. Chem. Soc. Jpn.*, 1975, **48**, 1191-1195.
55. A. A. Lamola and G. S. Hammond, *J. Chem. Phys.*, 1965, **43**, 2129-2135.
56. Y. L. Chow and X. Y. Liu, *Can. J. Chem.*, 1991, **69**, 1261-1272.
57. F. M. Richards, *J. Mol. Biol.*, 1974, **82**, 1-14.
58. M. Rivera, M. Dommett, A. Sidat, W. Rahim and R. Crespo-Otero, *J. Comput. Chem.*, 2020, **41**, 1045-1058.
59. X. Cai, M. Sakamoto, M. Yamaji, M. Fujitsuka and T. Majima, *Chem. Eur. J.*, 2007, **13**, 3143-3149.
60. J. C. Scaiano, B. R. Arnold and W. G. McGimpsey, *J. Phys. Chem.*, 1994, **98**, 5431-5434.
61. S. A. Green, D. J. Simpson, G. Zhou, P. S. Ho and N. V. Blough, *J. Am. Chem. Soc.*, 1990, **112**, 7337-7346.
62. A. I. Privalova, J. P. Morozova, E. R. Kashapova and V. J. Artyukhov, *J. Appl. Spectrosc.*, 2011, **78**, 309-317.
63. A. Mohan, E. Sebastian, M. Gudem and M. Hariharan, *J. Phys. Chem. B*, 2020, **124**, 6867-6874.
64. D. Beljonne, Z. Shuai, G. Pourtois and J. L. Bredas, *J. Phys. Chem. A*, 2001, **105**, 3899-3907.
65. M. J. Frisch, G. W. Trucks, H. B. Schlegel, G. E. Scuseria, et al., Gaussian 16 Rev. C.01, Gaussian, Inc., Wallingford, CT, 2016.
66. M. J. G. Peach and D. J. Tozer, *J. Phys. Chem. A*, 2012, **116**, 9783-9789.



## Table of contents entry

View Article Online  
DOI: 10.1039/D0SC04646A

The number and position of halogen substituents in purely organic  $\pi$ - $\pi^*$  chromophores critically affect the efficiency of phosphorescence.

